

# Direct WIMP Detection

*Manuel Drees*

*Bonn University & Bethe Center for Theoretical Physics*



# Contents

## 1 WIMP Dark Matter

# Contents

1 WIMP Dark Matter

2 Learning from Direct Detection

# Contents

1 WIMP Dark Matter

2 Learning from Direct Detection

a) Detection Principle

# Contents

1 WIMP Dark Matter

2 Learning from Direct Detection

a) Detection Principle

b) Velocity Distribution

# Contents

## 1 WIMP Dark Matter

## 2 Learning from Direct Detection

- a) Detection Principle
- b) Velocity Distribution
- c) WIMP Mass

# Contents

## 1 WIMP Dark Matter

## 2 Learning from Direct Detection

- a) Detection Principle
- b) Velocity Distribution
- c) WIMP Mass
- d) Cross Section times Density

# Contents

## 1 WIMP Dark Matter

## 2 Learning from Direct Detection

a) Detection Principle

b) Velocity Distribution

c) WIMP Mass

d) Cross Section times Density

## 3 Learning about Direct Detection in SUSY



# Contents

## 1 WIMP Dark Matter

## 2 Learning from Direct Detection

- a) Detection Principle
- b) Velocity Distribution
- c) WIMP Mass
- d) Cross Section times Density

## 3 Learning about Direct Detection in SUSY

- a) Spin-Independent Cross Section

# Contents

## 1 WIMP Dark Matter

## 2 Learning from Direct Detection

- a) Detection Principle
- b) Velocity Distribution
- c) WIMP Mass
- d) Cross Section times Density

## 3 Learning about Direct Detection in SUSY

- a) Spin-Independent Cross Section
- b) Hadronic Uncertainty

# Contents

## 1 WIMP Dark Matter

## 2 Learning from Direct Detection

- a) Detection Principle
- b) Velocity Distribution
- c) WIMP Mass
- d) Cross Section times Density

## 3 Learning about Direct Detection in SUSY

- a) Spin-Independent Cross Section
- b) Hadronic Uncertainty
- c) Sensitivity to SUSY Parameters

# Contents

## 1 WIMP Dark Matter

## 2 Learning from Direct Detection

- a) Detection Principle
- b) Velocity Distribution
- c) WIMP Mass
- d) Cross Section times Density

## 3 Learning about Direct Detection in SUSY

- a) Spin-Independent Cross Section
- b) Hadronic Uncertainty
- c) Sensitivity to SUSY Parameters

## 4 Summary

# Introduction: WIMPs as Dark Matter

Several observations indicate existence of non-luminous  
**Dark Matter (DM)** (more exactly: missing force)

# Introduction: WIMPs as Dark Matter

Several observations indicate existence of non-luminous **Dark Matter (DM)** (more exactly: missing force)

- **Galactic rotation curves** imply  $\Omega_{\text{DM}}h^2 \geq 0.05$ .

$\Omega$ : Mass density in units of critical density;  $\Omega = 1$  means flat Universe.

$h$ : Scaled Hubble constant. Observation:  $h = 0.72 \pm 0.07$  (?)

# Introduction: WIMPs as Dark Matter

Several observations indicate existence of non-luminous **Dark Matter (DM)** (more exactly: missing force)

- **Galactic rotation curves** imply  $\Omega_{\text{DM}}h^2 \geq 0.05$ .

$\Omega$ : Mass density in units of critical density;  $\Omega = 1$  means flat Universe.

$h$ : Scaled Hubble constant. Observation:  $h = 0.72 \pm 0.07$  (?)

- Models of structure formation,  $X$  ray temperature of clusters of galaxies, ...

# Introduction: WIMPs as Dark Matter

Several observations indicate existence of non-luminous **Dark Matter (DM)** (more exactly: missing force)

- **Galactic rotation curves** imply  $\Omega_{\text{DM}}h^2 \geq 0.05$ .

$\Omega$ : Mass density in units of critical density;  $\Omega = 1$  means flat Universe.

$h$ : Scaled Hubble constant. Observation:  $h = 0.72 \pm 0.07$  (?)

- Models of structure formation,  $X$  ray temperature of clusters of galaxies, ...
- **Cosmic Microwave Background anisotropies (WMAP)**  
imply  $\Omega_{\text{DM}}h^2 = 0.105^{+0.007}_{-0.013}$  Spergel et al., astro-ph/0603449



# Weakly Interacting Massive Particles (WIMPs)

- Exist in well-motivated extensions of the SM: SUSY, (Little Higgs with  $T$ -Parity), ((Universal Extra Dimension))

# Weakly Interacting Massive Particles (WIMPs)

- Exist in well-motivated extensions of the SM: SUSY, (Little Higgs with  $T$ -Parity), ((Universal Extra Dimension))
- Can also (trivially) write down “tailor-made” WIMP models

# Weakly Interacting Massive Particles (WIMPs)

- Exist in well-motivated extensions of the SM: SUSY, (Little Higgs with  $T$ -Parity), ((Universal Extra Dimension))
- Can also (trivially) write down “tailor-made” WIMP models
- In standard cosmology, roughly weak cross section automatically gives roughly right relic density for thermal WIMPs! (On logarithmic scale)

# Weakly Interacting Massive Particles (WIMPs)

- Exist in well-motivated extensions of the SM: SUSY, (Little Higgs with  $T$ -Parity), ((Universal Extra Dimension))
- Can also (trivially) write down “tailor-made” WIMP models
- In standard cosmology, roughly weak cross section automatically gives roughly right relic density for thermal WIMPs! (On logarithmic scale)
- Roughly weak interactions may allow both indirect and *direct* detection of WIMPs

# Probing WIMPs

Detection of WIMP annihilation products (“indirect detection”) suffers from uncertainties in

- Backgrounds

# Probing WIMPs

Detection of WIMP annihilation products (“indirect detection”) suffers from uncertainties in

- Backgrounds
- Propagation, esp. for charged particles ( $e^+$ ,  $\bar{p}$ ,  $\bar{d}$ )

# Probing WIMPs

Detection of WIMP annihilation products (“indirect detection”) suffers from uncertainties in

- Backgrounds
- Propagation, esp. for charged particles ( $e^+$ ,  $\bar{p}$ ,  $\bar{d}$ )

WIMP production at colliders: Can't be sure that

- WIMP is stable on cosmological time scales

# Probing WIMPs

Detection of WIMP annihilation products (“indirect detection”) suffers from uncertainties in

- Backgrounds
- Propagation, esp. for charged particles ( $e^+$ ,  $\bar{p}$ ,  $\bar{d}$ )

WIMP production at colliders: Can't be sure that

- WIMP is stable on cosmological time scales
- Cosmology is right



# Direct WIMP Detection

Look for elastic scattering of ambient WIMP off nucleus in detector; measure nuclear recoil energy.

# Direct WIMP Detection

Look for elastic scattering of ambient WIMP off nucleus in detector; measure nuclear recoil energy.

Direct WIMP detection is easiest convincing way to prove that WIMPs form DM!

# Direct WIMP Detection

Look for elastic scattering of ambient WIMP off nucleus in detector; measure nuclear recoil energy.

Direct WIMP detection is easiest convincing way to prove that WIMPs form DM! Other possibilities:  $\gamma$  line, high- $E$   $\nu$ 's from Sun: less "likely" to work.

# Direct WIMP Detection

Look for elastic scattering of ambient WIMP off nucleus in detector; measure nuclear recoil energy.

Direct WIMP detection is easiest convincing way to prove that WIMPs form DM! Other possibilities:  $\gamma$  line, high- $E$   $\nu$ 's from Sun: less "likely" to work.

Can also be interesting probe!

# Direct WIMP Detection: Formalism

Counting rate given by

$$\frac{dR}{dQ} = AF^2(Q) \int_{v_{\min}}^{v_{\max}} \frac{f_1(v)}{v} dv$$

$Q$ : recoil energy

$A = \rho\sigma_0 / (2m_\chi m_r) = \text{const.}$ : encodes particle physics

$F(Q)$ : nuclear form factor

$v$ : WIMP velocity in lab frame

$v_{\min}^2 = m_N Q / (2m_r^2)$  ( $m_r$ : reduced mass)

$v_{\max}$ : Maximal velocity of WIMPs bound to galaxy

$f_1(v)$ : normalized one-dimensional WIMP velocity distribution

Note:  $Q^2 \propto v^2(1 - \cos\theta^*) \Rightarrow \frac{d\sigma}{dQ} \propto \frac{1}{v^2} \frac{d\sigma}{d\cos\theta^*}$ .

# Direct WIMP Detection: Formalism

Counting rate given by

$$\frac{dR}{dQ} = AF^2(Q) \int_{v_{\min}}^{v_{\max}} \frac{f_1(v)}{v} dv$$

$Q$ : recoil energy

$A = \rho\sigma_0 / (2m_\chi m_r) = \text{const.}$ : encodes particle physics

$F(Q)$ : nuclear form factor

$v$ : WIMP velocity in lab frame

$v_{\min}^2 = m_N Q / (2m_r^2)$  ( $m_r$ : reduced mass)

$v_{\max}$ : Maximal velocity of WIMPs bound to galaxy

$f_1(v)$ : normalized one-dimensional WIMP velocity distribution

Note:  $Q^2 \propto v^2(1 - \cos\theta^*) \Rightarrow \frac{d\sigma}{dQ} \propto \frac{1}{v^2} \frac{d\sigma}{d\cos\theta^*}$ .

Can invert this relation to measure  $f_1(v)$ !

# Direct reconstruction of $f_1$

MD & C.L. Shan, astro-ph/0703651

$$f_1(v) = \mathcal{N} \left\{ -2Q \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \frac{dR}{dQ} \right] \right\}_{Q=2m_r^2 v^2 / m_N}$$

# Direct reconstruction of $f_1$

MD & C.L. Shan, astro-ph/0703651

$$f_1(v) = \mathcal{N} \left\{ -2Q \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \frac{dR}{dQ} \right] \right\}_{Q=2m_r^2 v^2 / m_N}$$

$\mathcal{N}$ : Normalization ( $\int_0^\infty f_1(v) dv = 1$ ).



# Direct reconstruction of $f_1$

MD & C.L. Shan, astro-ph/0703651

$$f_1(v) = \mathcal{N} \left\{ -2Q \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \frac{dR}{dQ} \right] \right\}_{Q=2m_r^2 v^2 / m_N}$$

$\mathcal{N}$ : Normalization ( $\int_0^\infty f_1(v) dv = 1$ ).

Need to know form factor  $\implies$  stick to spin-independent scattering.

# Direct reconstruction of $f_1$

MD & C.L. Shan, astro-ph/0703651

$$f_1(v) = \mathcal{N} \left\{ -2Q \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \frac{dR}{dQ} \right] \right\}_{Q=2m_r^2 v^2 / m_N}$$

$\mathcal{N}$ : Normalization ( $\int_0^\infty f_1(v) dv = 1$ ).

Need to know form factor  $\implies$  stick to spin-independent scattering.

Need to know  $m_\chi$ , but do *not* need  $\sigma_0, \rho$ .

# Direct reconstruction of $f_1$

MD & C.L. Shan, astro-ph/0703651

$$f_1(v) = \mathcal{N} \left\{ -2Q \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \frac{dR}{dQ} \right] \right\}_{Q=2m_r^2 v^2 / m_N}$$

$\mathcal{N}$ : Normalization ( $\int_0^\infty f_1(v) dv = 1$ ).

Need to know form factor  $\implies$  stick to spin-independent scattering.

Need to know  $m_\chi$ , but do *not* need  $\sigma_0, \rho$ .

Need to know *slope* of recoil spectrum!

# Direct reconstruction of $f_1$

MD & C.L. Shan, astro-ph/0703651

$$f_1(v) = \mathcal{N} \left\{ -2Q \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \frac{dR}{dQ} \right] \right\}_{Q=2m_\chi^2 v^2 / m_N}$$

$\mathcal{N}$ : Normalization ( $\int_0^\infty f_1(v) dv = 1$ ).

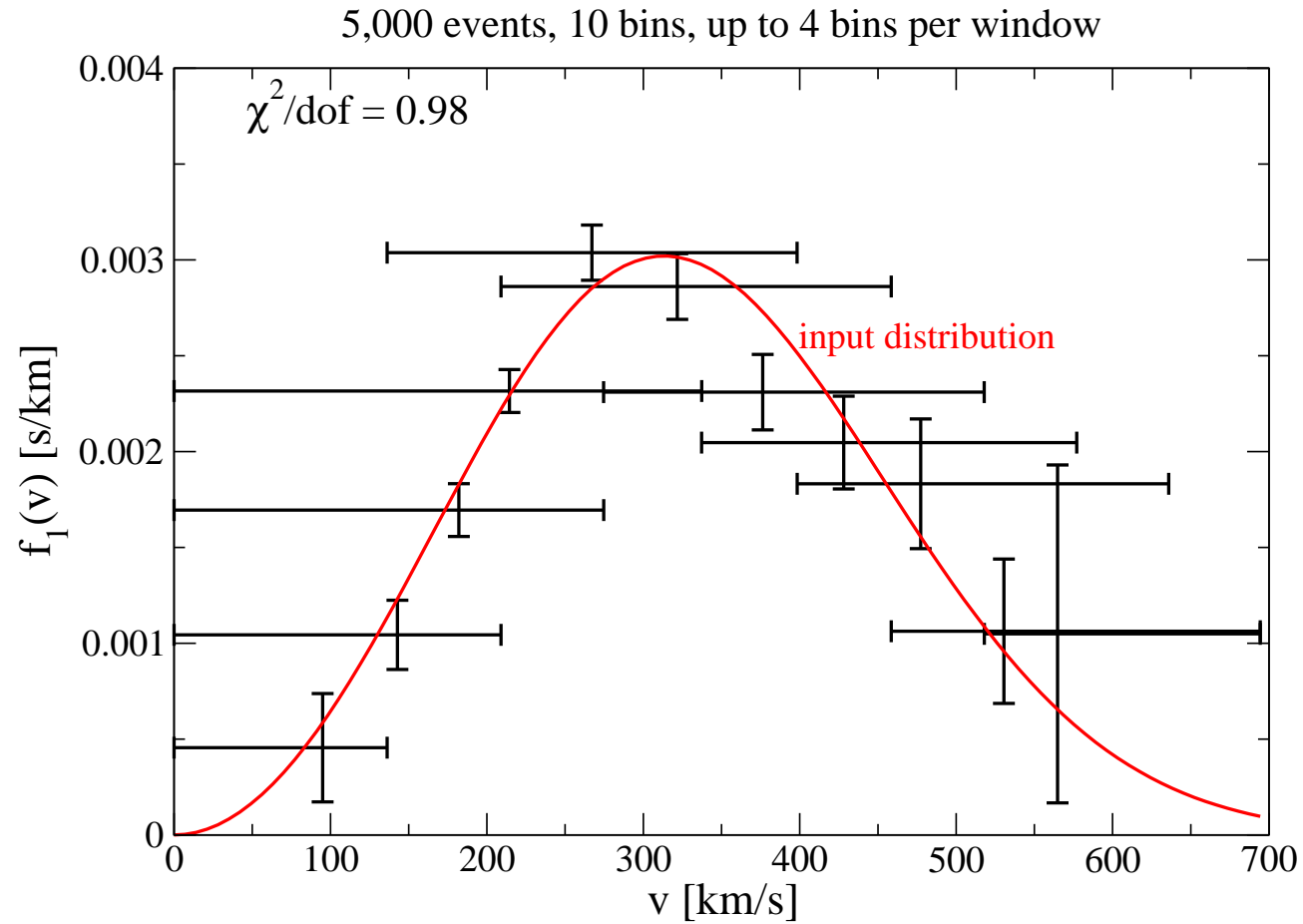
Need to know form factor  $\implies$  stick to spin-independent scattering.

Need to know  $m_\chi$ , but do *not* need  $\sigma_0, \rho$ .

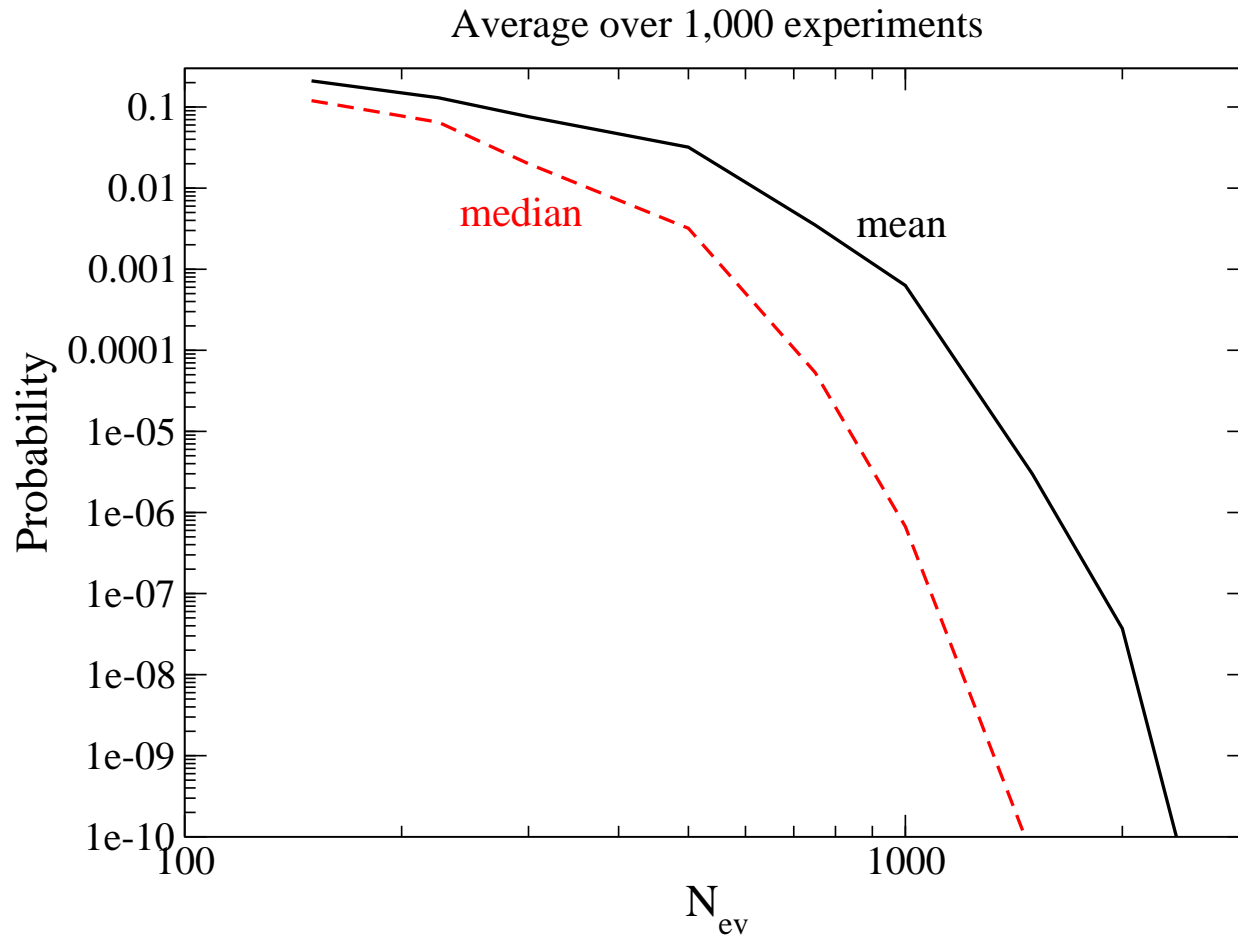
Need to know *slope* of recoil spectrum!

$dR/dQ$  is approximately exponential: better work with logarithmic slope: from  $\langle Q \rangle$  in bin!

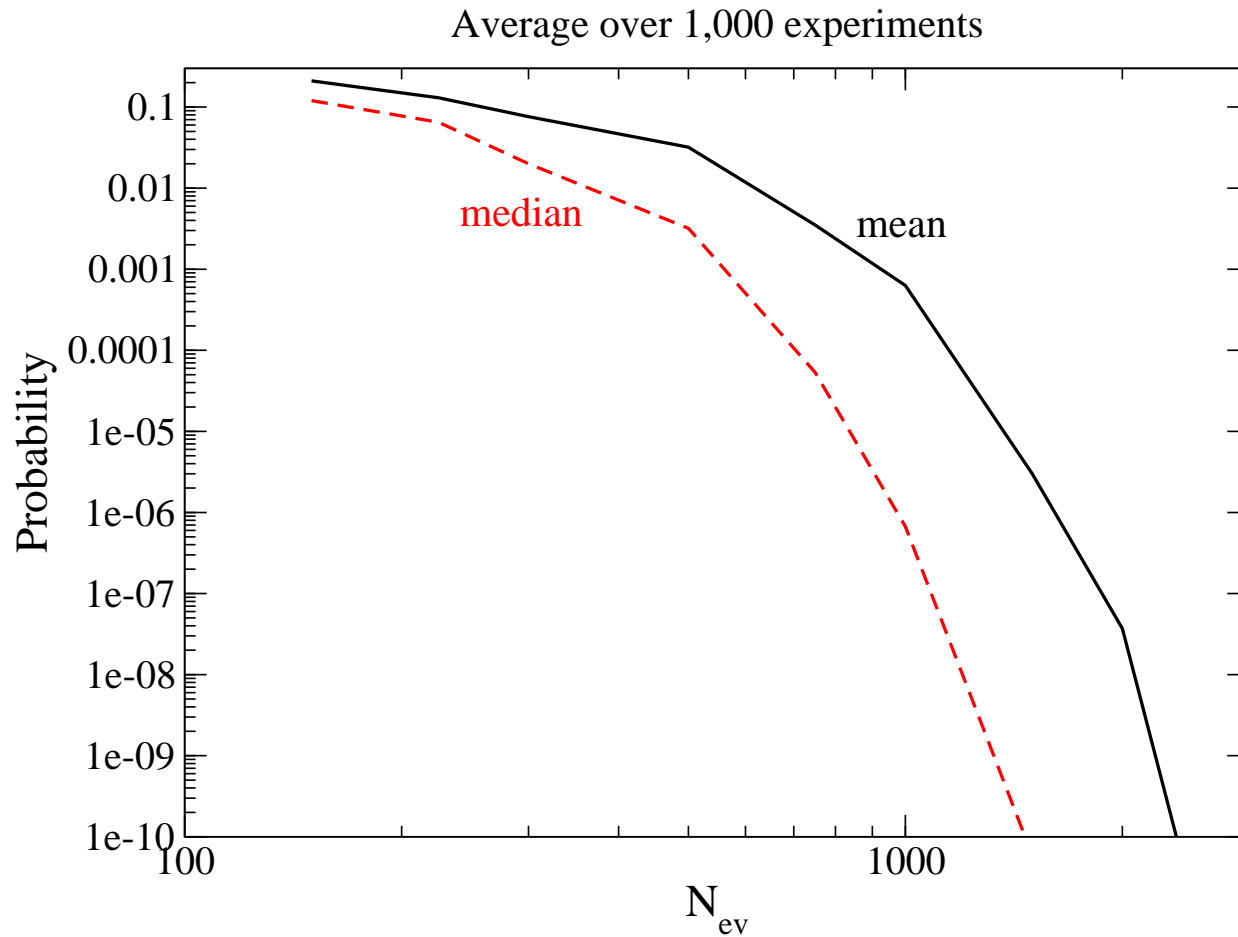
# Recoil spectrum: prediction and simulated measurement



# Statistical exclusion of constant $f_1$



# Statistical exclusion of constant $f_1$



Need several hundred events to begin direct reconstruction!

# Determining moments of $f_1$

$$\langle v^n \rangle \equiv \int_0^\infty v^n f_1(v) dv$$



# Determining moments of $f_1$

$$\begin{aligned}\langle v^n \rangle &\equiv \int_0^\infty v^n f_1(v) dv \\ &\propto \int_0^\infty Q^{(n-1)/2} \frac{1}{F^2(Q)} \frac{dR}{dQ} dQ\end{aligned}$$

# Determining moments of $f_1$

$$\begin{aligned}\langle v^n \rangle &\equiv \int_0^\infty v^n f_1(v) dv \\ &\propto \int_0^\infty Q^{(n-1)/2} \frac{1}{F^2(Q)} \frac{dR}{dQ} dQ \\ &\rightarrow \sum_{\text{events } a} \frac{Q_a^{(n-1)/2}}{F^2(Q_a)}\end{aligned}$$

# Determining moments of $f_1$

$$\begin{aligned}\langle v^n \rangle &\equiv \int_0^\infty v^n f_1(v) dv \\ &\propto \int_0^\infty Q^{(n-1)/2} \frac{1}{F^2(Q)} \frac{dR}{dQ} dQ \\ &\rightarrow \sum_{\text{events } a} \frac{Q_a^{(n-1)/2}}{F^2(Q_a)}\end{aligned}$$

Can incorporate finite energy (hence velocity) threshold

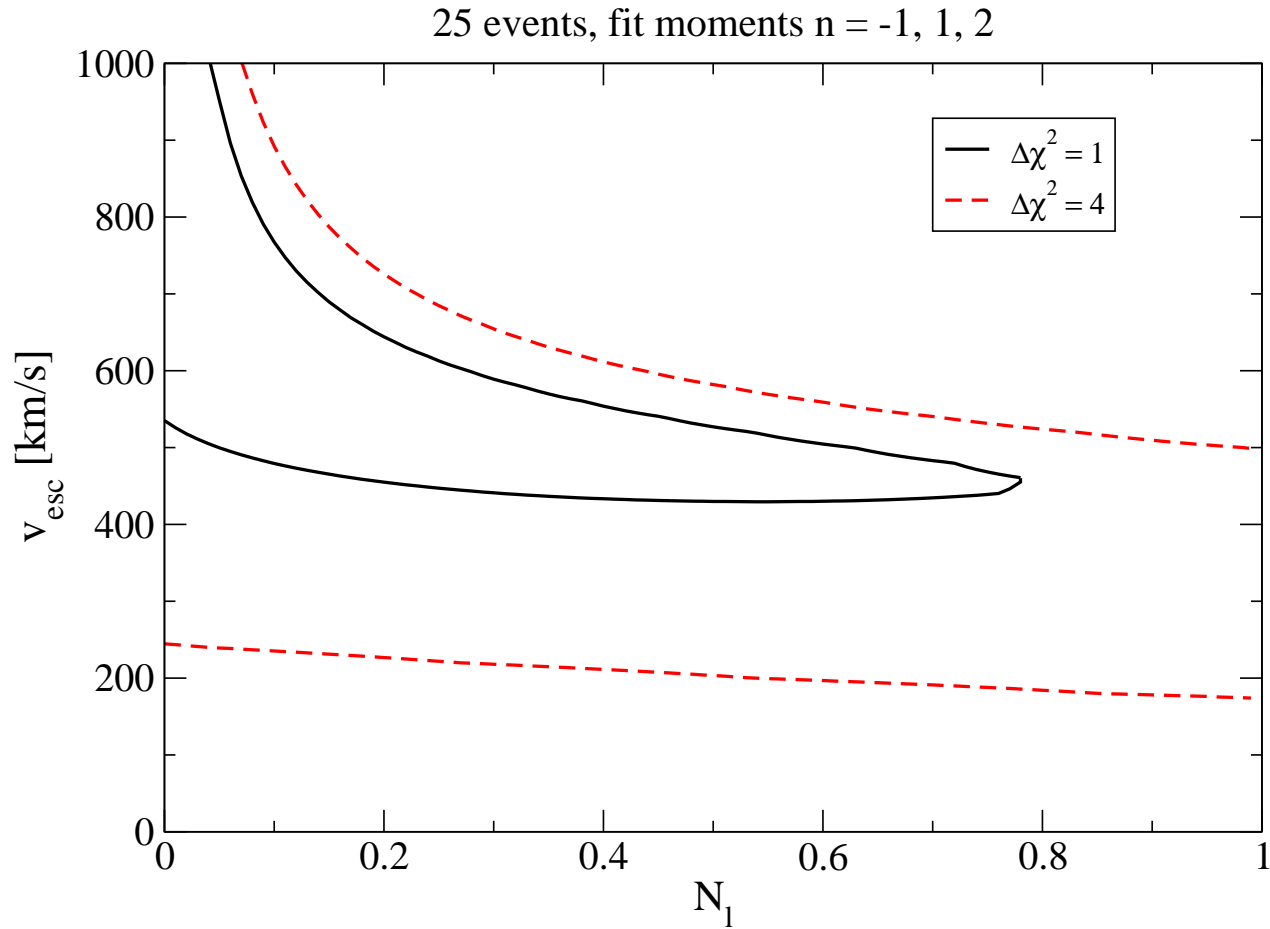
# Determining moments of $f_1$

$$\begin{aligned}\langle v^n \rangle &\equiv \int_0^\infty v^n f_1(v) dv \\ &\propto \int_0^\infty Q^{(n-1)/2} \frac{1}{F^2(Q)} \frac{dR}{dQ} dQ \\ &\rightarrow \sum_{\text{events } a} \frac{Q_a^{(n-1)/2}}{F^2(Q_a)}\end{aligned}$$

Can incorporate finite energy (hence velocity) threshold

Moments are strongly correlated!

# Constraining a “late infall” component



# Determining the WIMP mass

MD & C.L. Shan, arXiv:0803447 (hep-ph)

- Method described above yields normalized  $f_1(v)$  for *any* assumed  $m_\chi$

# Determining the WIMP mass

MD & C.L. Shan, arXiv:0803447 (hep-ph)

- Method described above yields normalized  $f_1(v)$  for *any* assumed  $m_\chi$
- $\Rightarrow$  *cannot* determine  $m_\chi$  from single recoil spectrum, unless  $f_1(v)$  is (assumed to be) known

# Determining the WIMP mass

MD & C.L. Shan, arXiv:0803447 (hep-ph)

- Method described above yields normalized  $f_1(v)$  for *any* assumed  $m_\chi$
- $\Rightarrow$  *cannot* determine  $m_\chi$  from single recoil spectrum, unless  $f_1(v)$  is (assumed to be) known
- *Can* determine  $m_\chi$  model-independently from two (or more) measurements, by demanding that they yield the same (moments of)  $f_1$ !



# Determining the WIMP mass

MD & C.L. Shan, arXiv:0803447 (hep-ph)

- Method described above yields normalized  $f_1(v)$  for *any* assumed  $m_\chi$
- $\Rightarrow$  *cannot* determine  $m_\chi$  from single recoil spectrum, unless  $f_1(v)$  is (assumed to be) known
- *Can* determine  $m_\chi$  model-independently from two (or more) measurements, by demanding that they yield the same (moments of)  $f_1$ !
- *Can* also get  $m_\chi$  from comparison of event rates, assuming equal cross section on neutrons and protons.

# Systematic errors

- Equality of moments of  $f_1$  holds only if integrals run over identical ranges of  $v$ , e.g.  $v_{\min} = 0$ ,  $v_{\max} = \infty$ .

# Systematic errors

- Equality of moments of  $f_1$  holds only if integrals run over identical ranges of  $v$ , e.g.  $v_{\min} = 0$ ,  $v_{\max} = \infty$ .
- Real experiments have finite acceptance windows for  $Q$ , and hence for  $v$

# Systematic errors

- Equality of moments of  $f_1$  holds only if integrals run over identical ranges of  $v$ , e.g.  $v_{\min} = 0$ ,  $v_{\max} = \infty$ .
- Real experiments have finite acceptance windows for  $Q$ , and hence for  $v$
- Ensuring  $v_{\min,X} = v_{\min,Y}$  and  $v_{\max,X} = v_{\max,Y}$  only possible if  $m_\chi$  is known

# Systematic errors

- Equality of moments of  $f_1$  holds only if integrals run over identical ranges of  $v$ , e.g.  $v_{\min} = 0$ ,  $v_{\max} = \infty$ .
- Real experiments have finite acceptance windows for  $Q$ , and hence for  $v$
- Ensuring  $v_{\min,X} = v_{\min,Y}$  and  $v_{\max,X} = v_{\max,Y}$  only possible if  $m_\chi$  is known
- For  $v_{\min}$ : Systematic effect not very large if  $m_\chi \gtrsim 20$  GeV,  $Q_{\min} \lesssim 3$  keV,  $Q_{\min,X} = Q_{\min,Y}$  terms included in  $I_n$ .

# Systematic errors

- Equality of moments of  $f_1$  holds only if integrals run over identical ranges of  $v$ , e.g.  $v_{\min} = 0$ ,  $v_{\max} = \infty$ .
- Real experiments have finite acceptance windows for  $Q$ , and hence for  $v$
- Ensuring  $v_{\min,X} = v_{\min,Y}$  and  $v_{\max,X} = v_{\max,Y}$  only possible if  $m_\chi$  is known
- For  $v_{\min}$ : Systematic effect not very large if  $m_\chi \gtrsim 20$  GeV,  $Q_{\min} \lesssim 3$  keV,  $Q_{\min,X} = Q_{\min,Y}$  terms included in  $I_n$ .
- Use  $Q_{\min} = 0$  from now on.

# Effect of finite $Q_{\max}$

- (Higher) moments are very sensitive to high- $Q$  region, even to region with  $\langle N_{\text{ev}} \rangle < 1$

# Effect of finite $Q_{\max}$

- (Higher) moments are very sensitive to high- $Q$  region, even to region with  $\langle N_{\text{ev}} \rangle < 1$
- Imposing finite  $Q_{\max}$  can alleviate this problem,



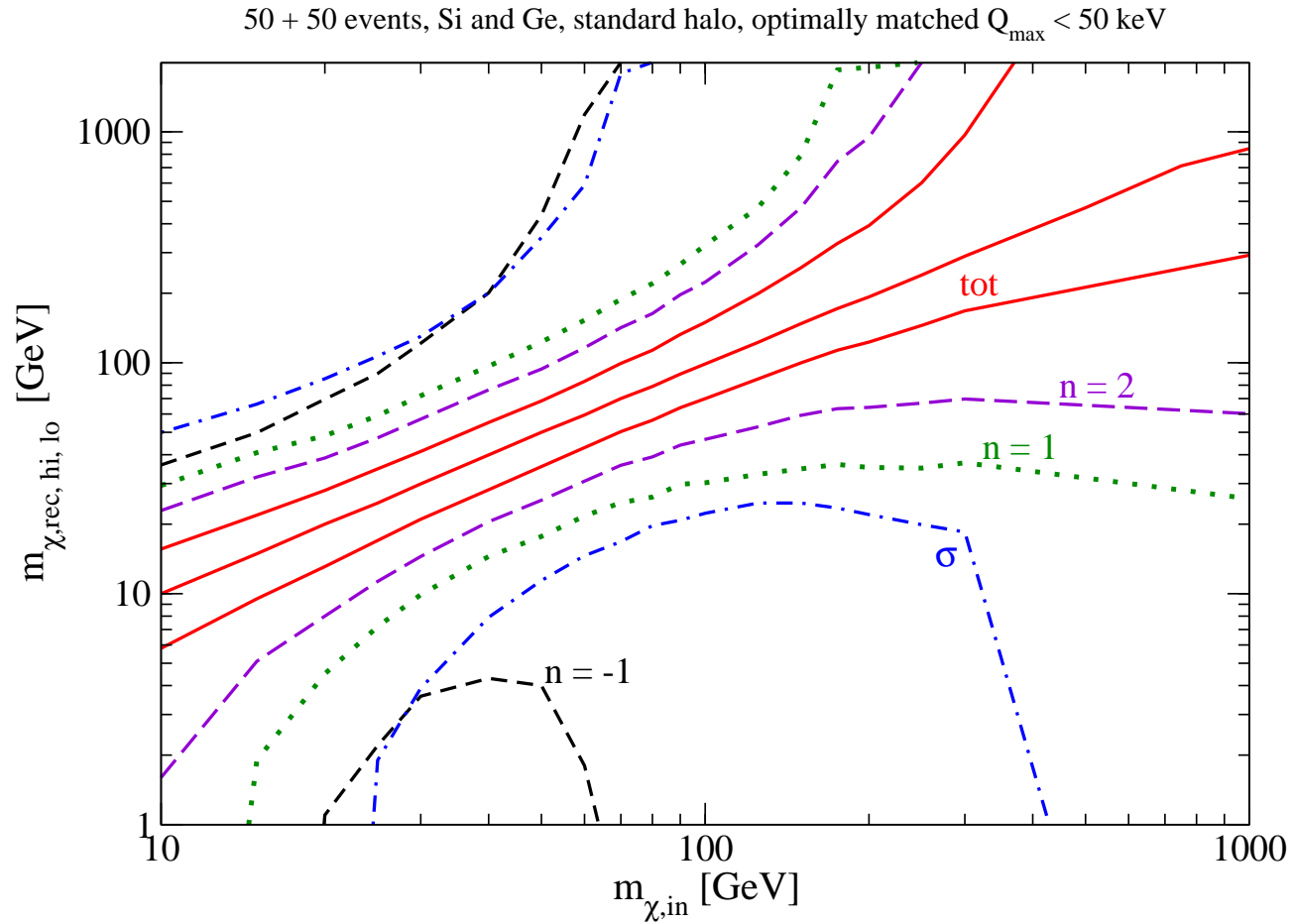
# Effect of finite $Q_{\max}$

- (Higher) moments are very sensitive to high- $Q$  region, even to region with  $\langle N_{\text{ev}} \rangle < 1$
- Imposing finite  $Q_{\max}$  can alleviate this problem,
- but introduces systematic error unless  $Q_{\max}$  values of two targets are matched; matching depends on  $m_\chi$ .

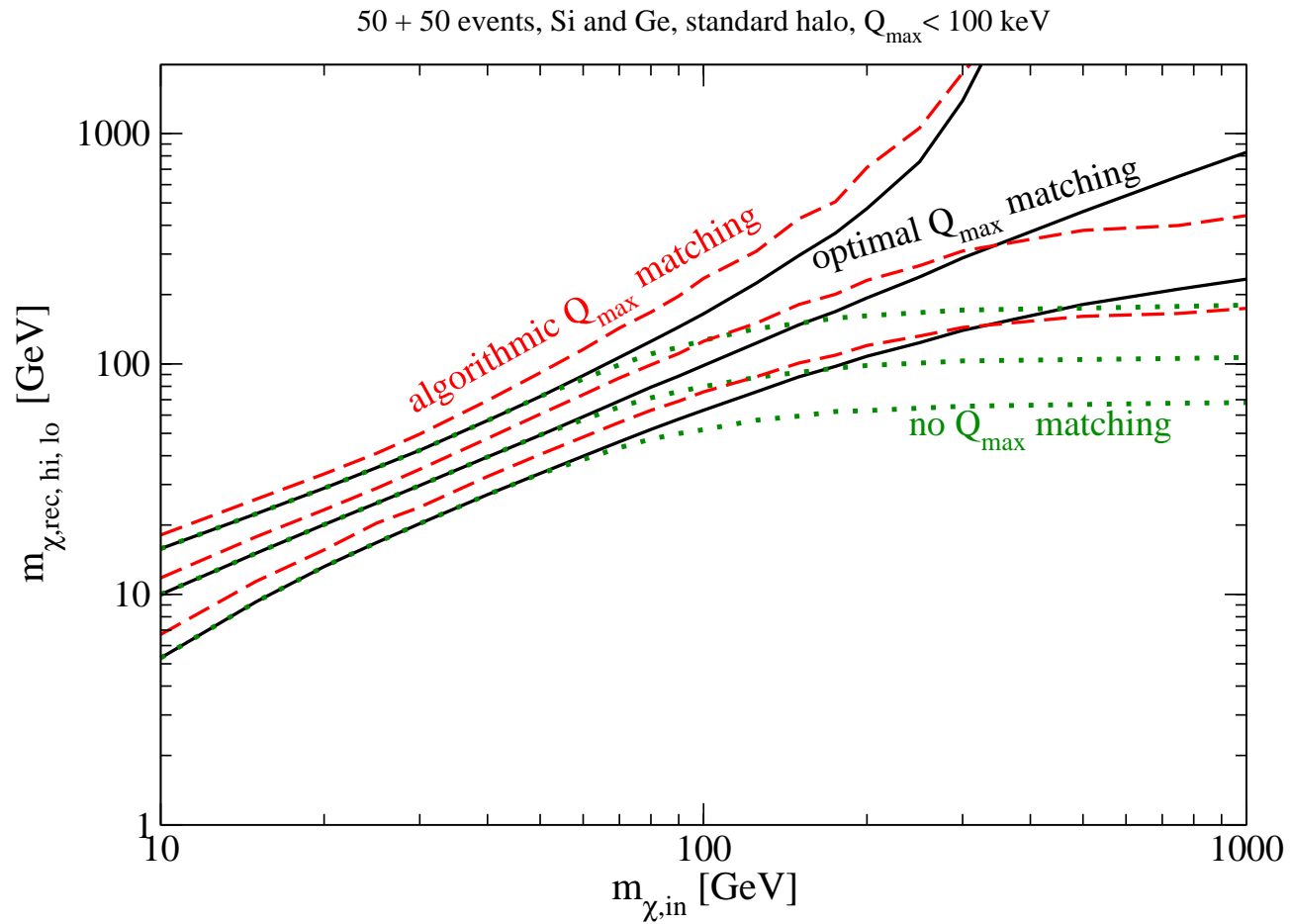
# Effect of finite $Q_{\max}$

- (Higher) moments are very sensitive to high- $Q$  region, even to region with  $\langle N_{\text{ev}} \rangle < 1$
- Imposing finite  $Q_{\max}$  can alleviate this problem,
- but introduces systematic error unless  $Q_{\max}$  values of two targets are matched; matching depends on  $m_\chi$ .
- Developed a method for this matching, based on  $\chi^2$  fit of several moments.

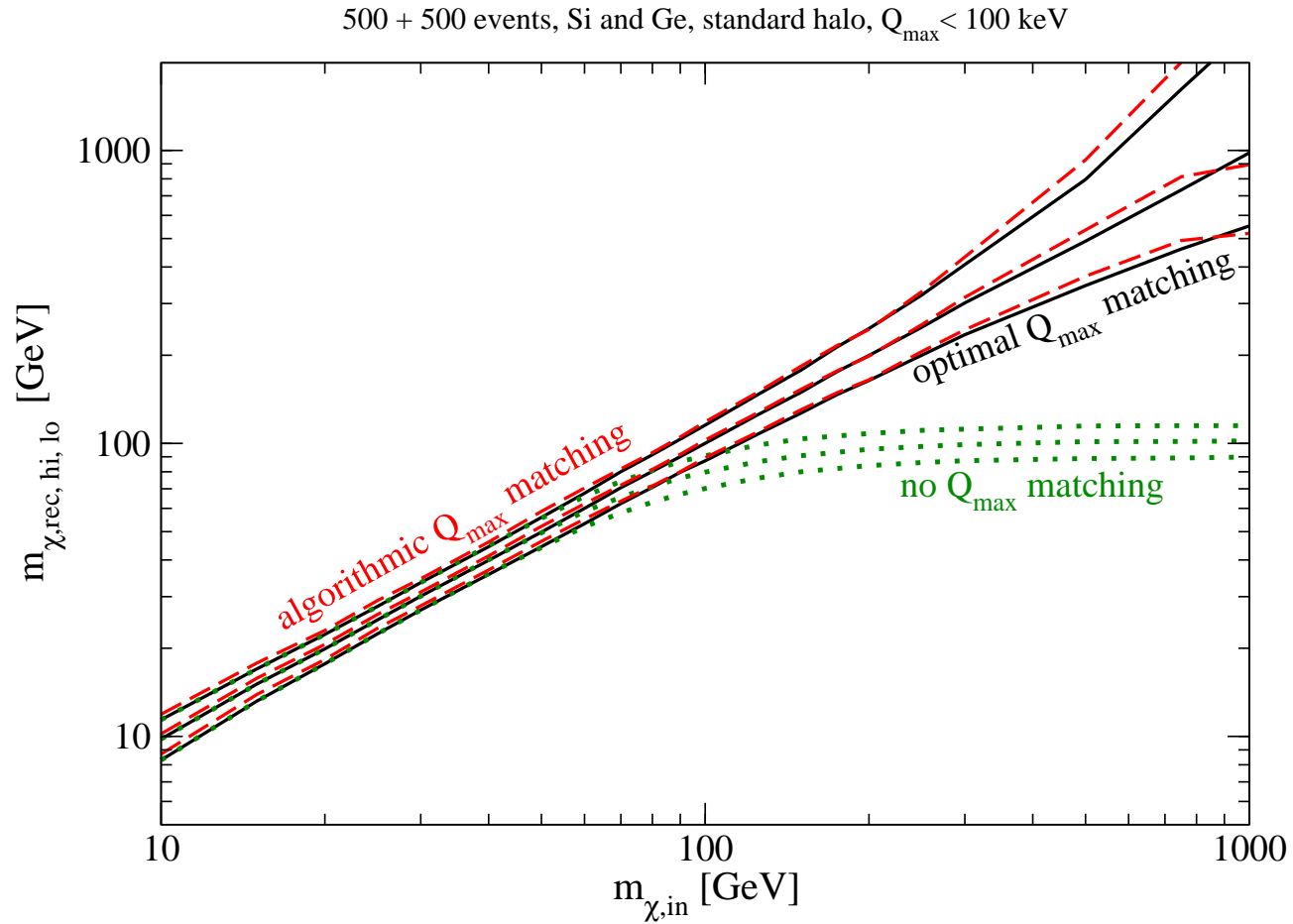
# Median reconstructed WIMP mass



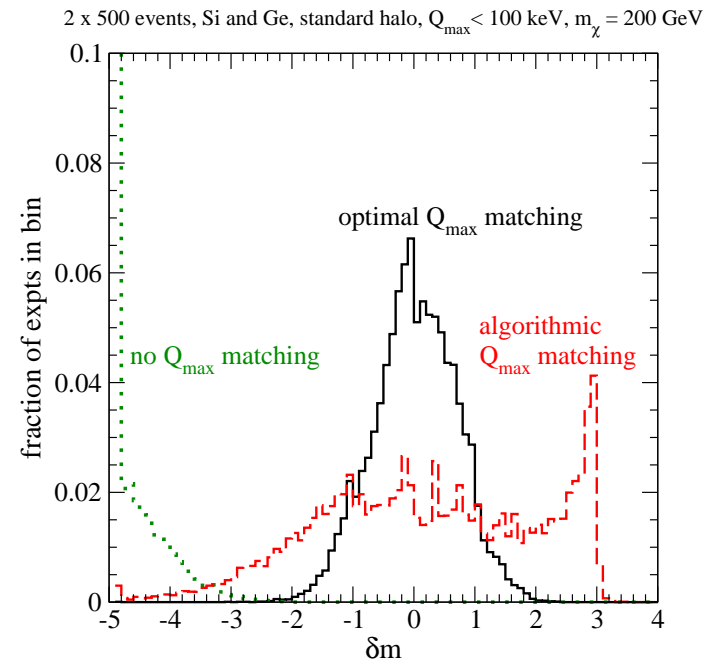
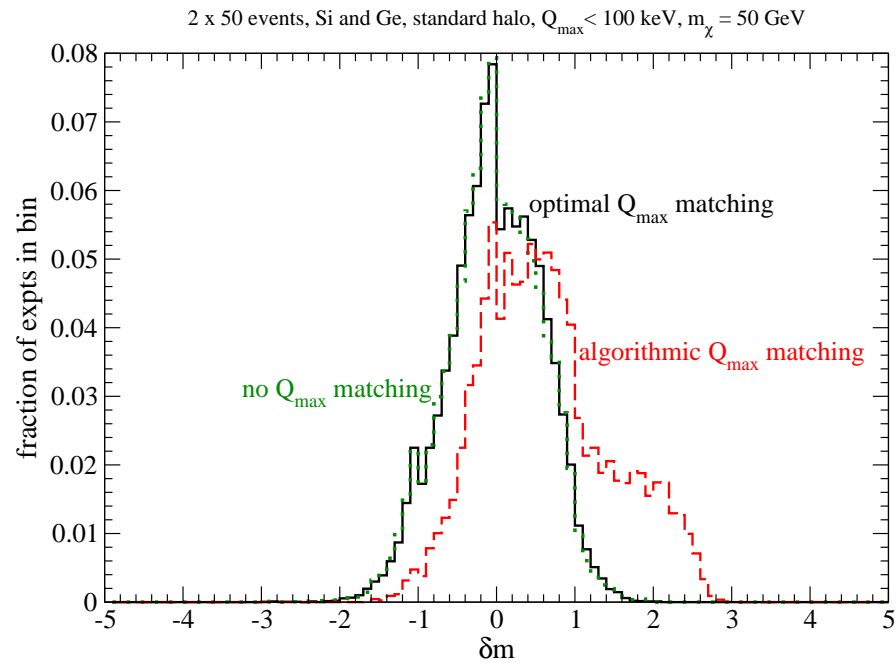
# Median reconstructed WIMP mass



# Median reconstructed WIMP mass



# Distribution of measurements



# WIMP Density times Cross Section

For spin-independent scattering:

$$\begin{aligned}\rho_\chi \sigma_{\chi p} &\propto \frac{r(Q_{\min})}{\langle v^{-1} \rangle} (m_\chi + m_N) \\ &\propto \left( \frac{2\sqrt{Q_{\min}} r(Q_{\min})}{F^2(Q_{\min})} + I_0 \right) (m_\chi + m_N). \quad (1)\end{aligned}$$

$$r(Q_{\min}) = \left. \frac{dR}{dQ} \right|_{Q=Q_{\min}}$$

First factor on r.h.s. in 2nd line comes from normalization of  $-1^{\text{st}}$  moment.

# WIMP Density times Cross Section

For spin-independent scattering:

$$\begin{aligned}\rho_\chi \sigma_{\chi p} &\propto \frac{r(Q_{\min})}{\langle v^{-1} \rangle} (m_\chi + m_N) \\ &\propto \left( \frac{2\sqrt{Q_{\min}} r(Q_{\min})}{F^2(Q_{\min})} + I_0 \right) (m_\chi + m_N). \quad (2)\end{aligned}$$

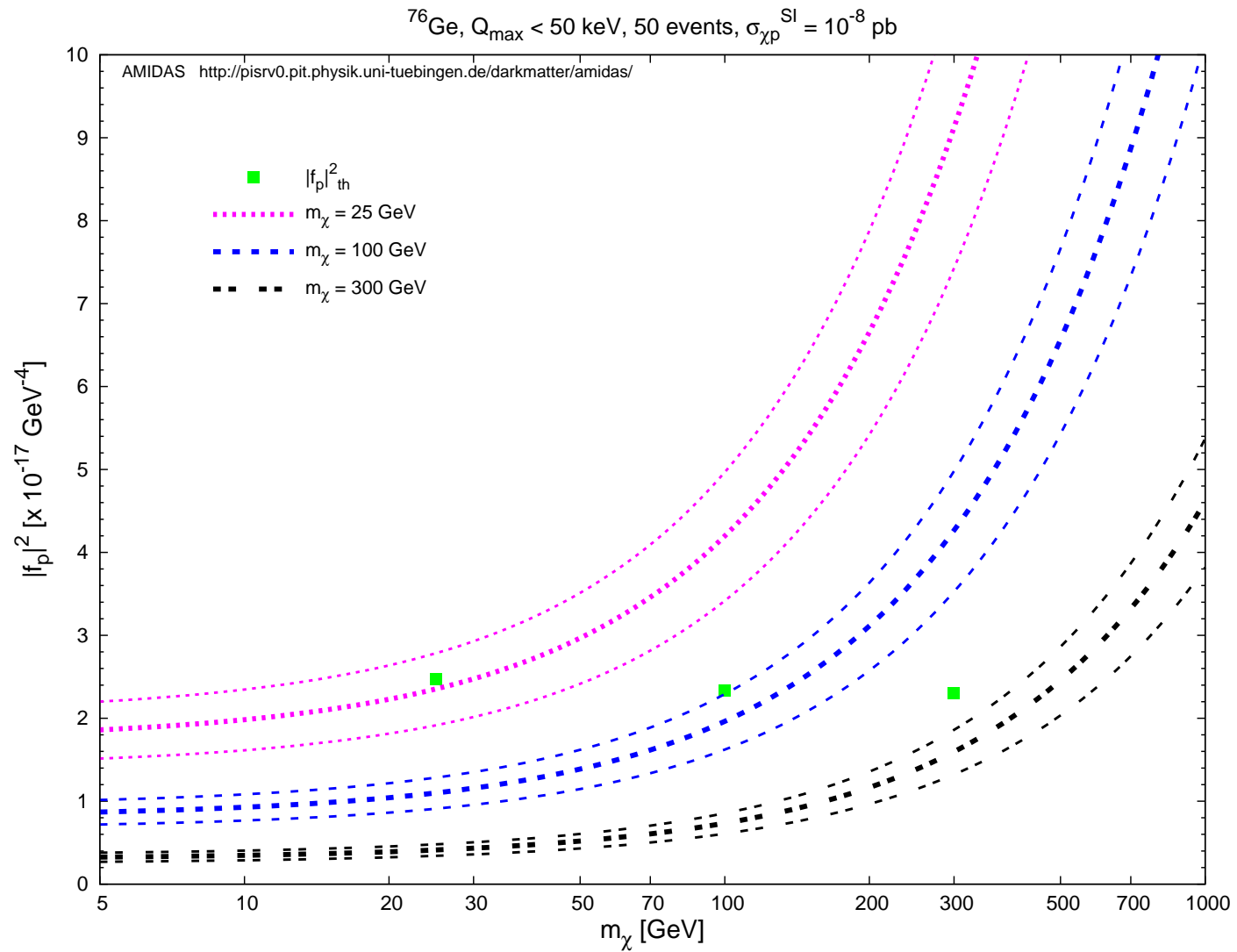
$$r(Q_{\min}) = \left. \frac{dR}{dQ} \right|_{Q=Q_{\min}}$$

First factor on r.h.s. in 2nd line comes from normalization of  $-1^{\text{st}}$  moment.

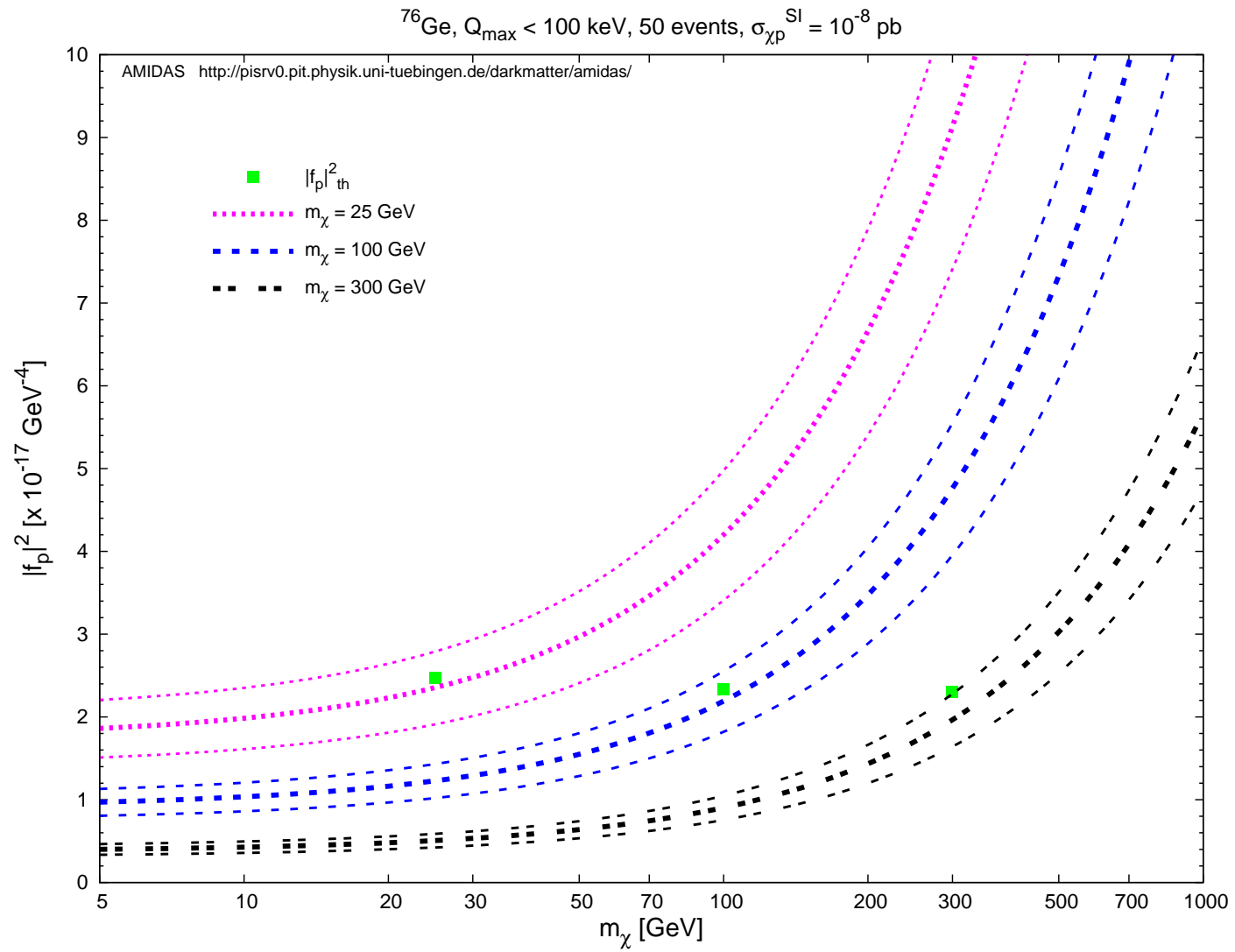
Can model-independently determine cross section times density from scattering data! MD & C.-L. Shan, to appear



# Results for $Q_{\max} = 50 \text{ keV}$



# Results for $Q_{\max} = 100 \text{ keV}$



# WIMP–Proton Scattering in SUSY

WIMP is lightest neutralino  $\tilde{\chi}_1^0$ .

# WIMP–Proton Scattering in SUSY

WIMP is lightest neutralino  $\tilde{\chi}_1^0$ .

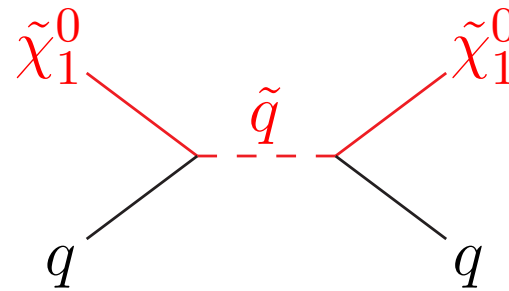
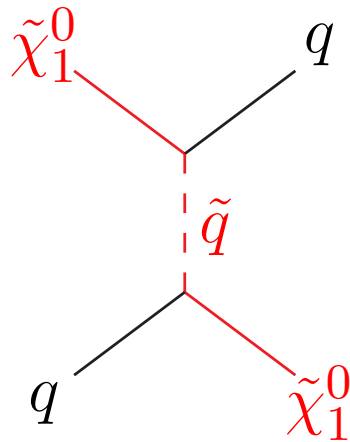
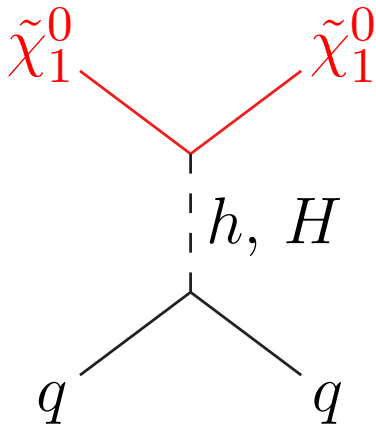
Stick to **spin-independent** contribution:  $\mathcal{L}_{\text{eff}} = f_p \bar{p} p \tilde{\chi}_1^0 \tilde{\chi}_1^0$

# WIMP–Proton Scattering in SUSY

WIMP is lightest neutralino  $\tilde{\chi}_1^0$ .

Stick to **spin-independent** contribution:  $\mathcal{L}_{\text{eff}} = f_p \bar{p} p \tilde{\chi}_1^0 \tilde{\chi}_1^0$

Come from spin-independent  $\tilde{\chi}_1^0 q$  interactions:

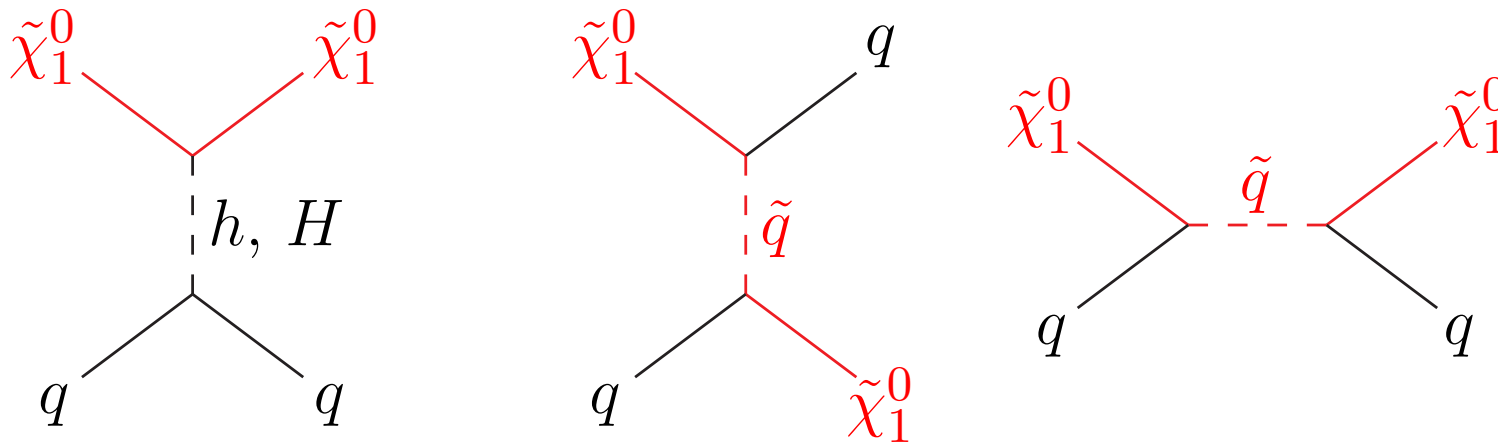


# WIMP–Proton Scattering in SUSY

WIMP is lightest neutralino  $\tilde{\chi}_1^0$ .

Stick to **spin-independent** contribution:  $\mathcal{L}_{\text{eff}} = f_p \bar{p} p \tilde{\chi}_1^0 \tilde{\chi}_1^0$

Come from spin-independent  $\tilde{\chi}_1^0 q$  interactions:



To  $\mathcal{O}(m_{\tilde{q}}^{-2})$ : Interaction  $\propto m_q$ ! From Higgs(ino) Yukawa,  $\tilde{q}_L - \tilde{q}_R$  mixing.

$\implies$  need matrix elements  $m_q \langle p | \bar{q} q | p \rangle$ !

## Matrix Elements $m_q \langle p | \bar{q}q | p \rangle$

- For heavy quarks,  $q = c, b, t$ : Calculate perturbatively via gluon loop. Shifman et al. 1977. Result is independent of  $m_q$ . Need some modification for  $\tilde{t}$  in loop. MD & Nojiri 1993

# Matrix Elements $m_q \langle p | \bar{q}q | p \rangle$

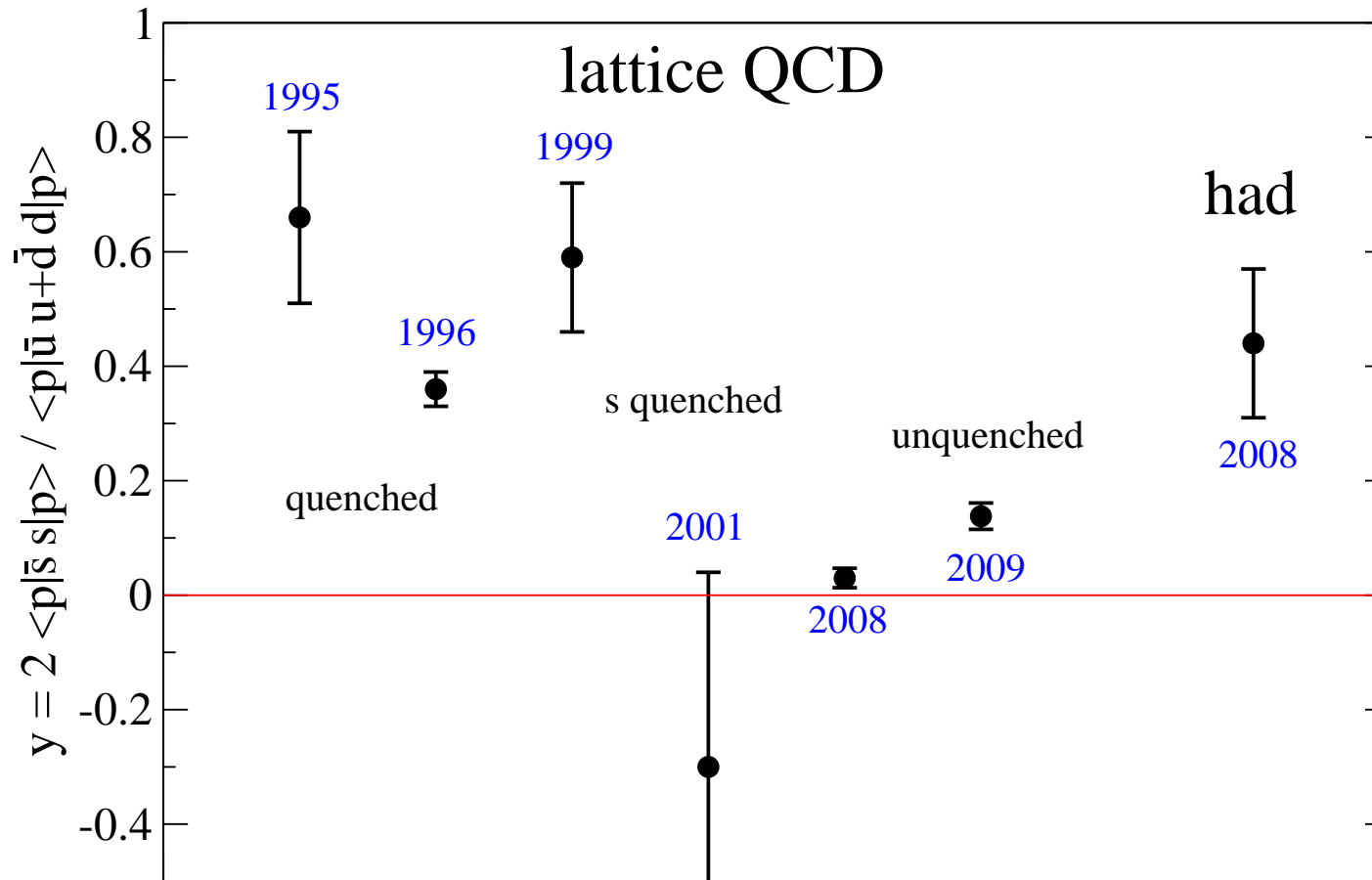
- For heavy quarks,  $q = c, b, t$ : Calculate perturbatively via gluon loop. Shifman et al. 1977. Result is independent of  $m_q$ . Need some modification for  $\tilde{t}$  in loop. MD & Nojiri 1993
- Need current quark masses  $\Rightarrow$  contributions from  $u, d$  are small  
 $\implies \sigma_{\tilde{\chi}_1^0 p} \simeq \sigma_{\tilde{\chi}_1^0 n}$  for spin-indep. contribution!



# Matrix Elements $m_q \langle p | \bar{q}q | p \rangle$

- For heavy quarks,  $q = c, b, t$ : Calculate perturbatively via gluon loop. Shifman et al. 1977. Result is independent of  $m_q$ . Need some modification for  $\tilde{t}$  in loop. MD & Nojiri 1993
- Need current quark masses  $\Rightarrow$  contributions from  $u, d$  are small  
 $\implies \sigma_{\tilde{\chi}_1^0 p} \simeq \sigma_{\tilde{\chi}_1^0 n}$  for spin-indep. contribution!
- Strange quark contribution important, but poorly known!

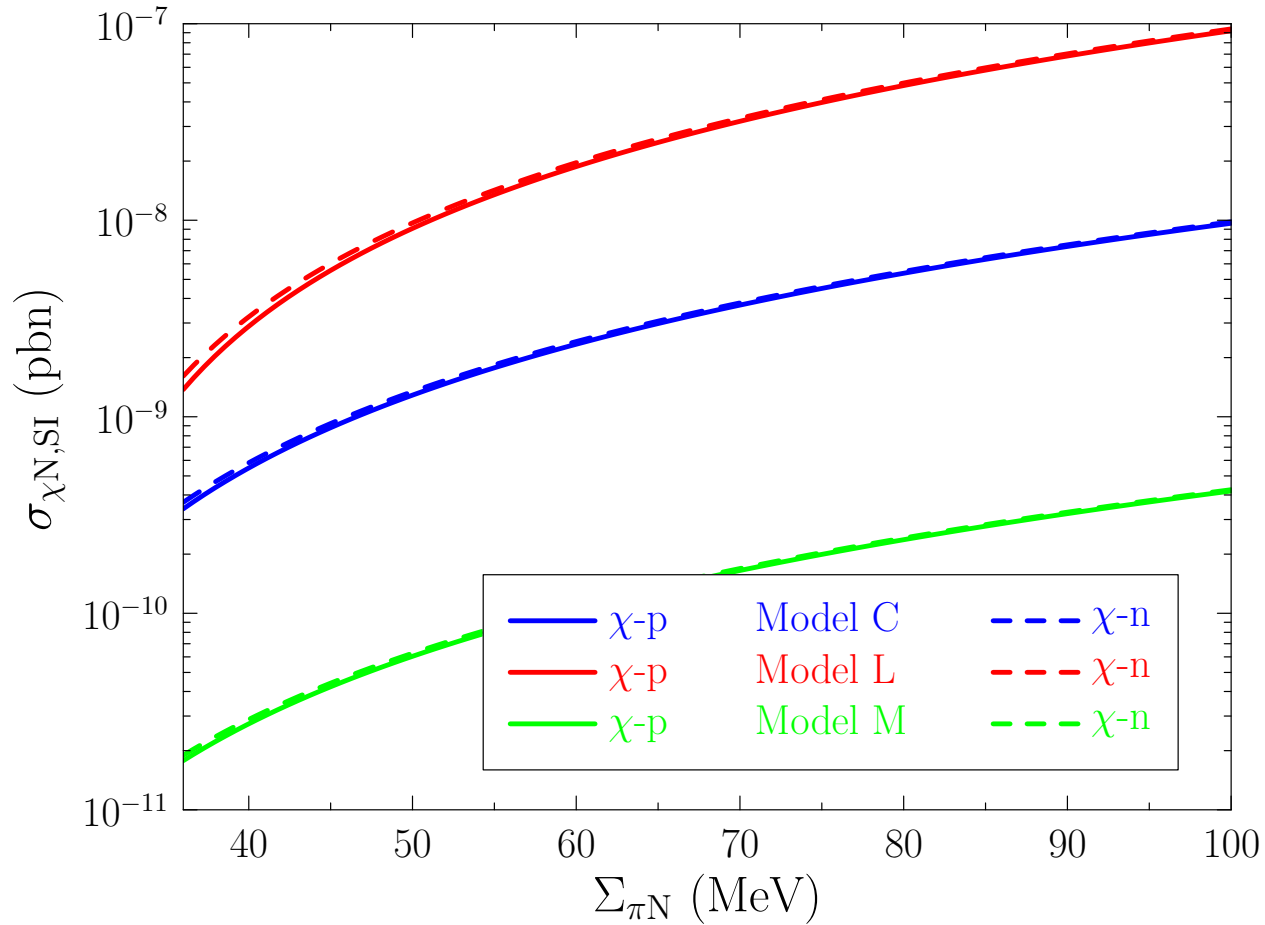
# Determinations of $\langle p | \bar{s}s | p \rangle$



Fukugita et al. (1995); Dong et al. (1996); Güsken et al. (1999); Michael et al. (2001); Ohki et al. (2008); Toussaint & Freeman (2009); Ellis et al. (2008)

# Effect of this uncertainty

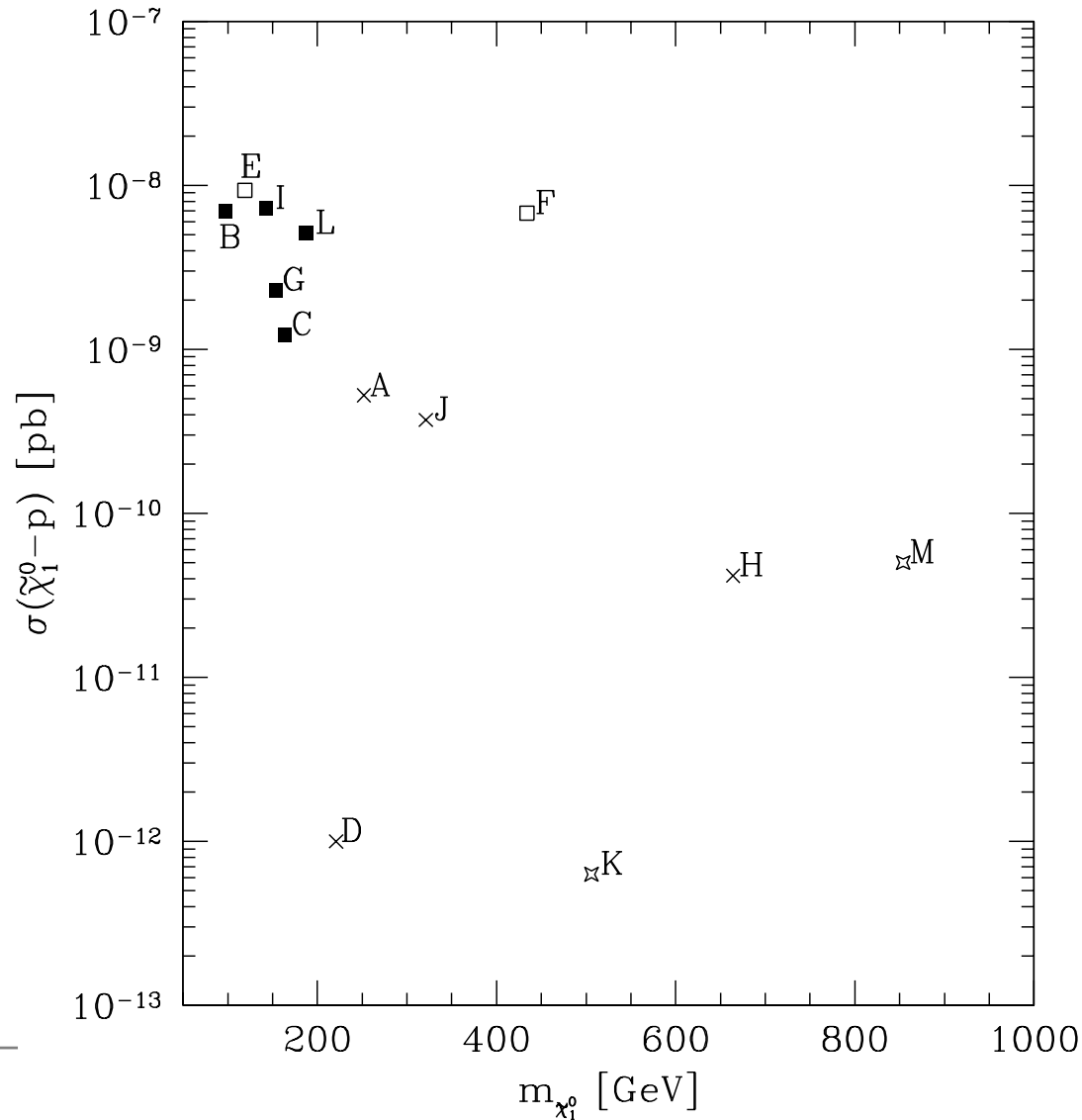
Ellis, Olive & Savage, arXiv:0801.3656



Larger  $\Sigma_{\pi N}$  implies larger  $\langle p | \bar{s}s | p \rangle$ .

# Survey of Benchmark Points

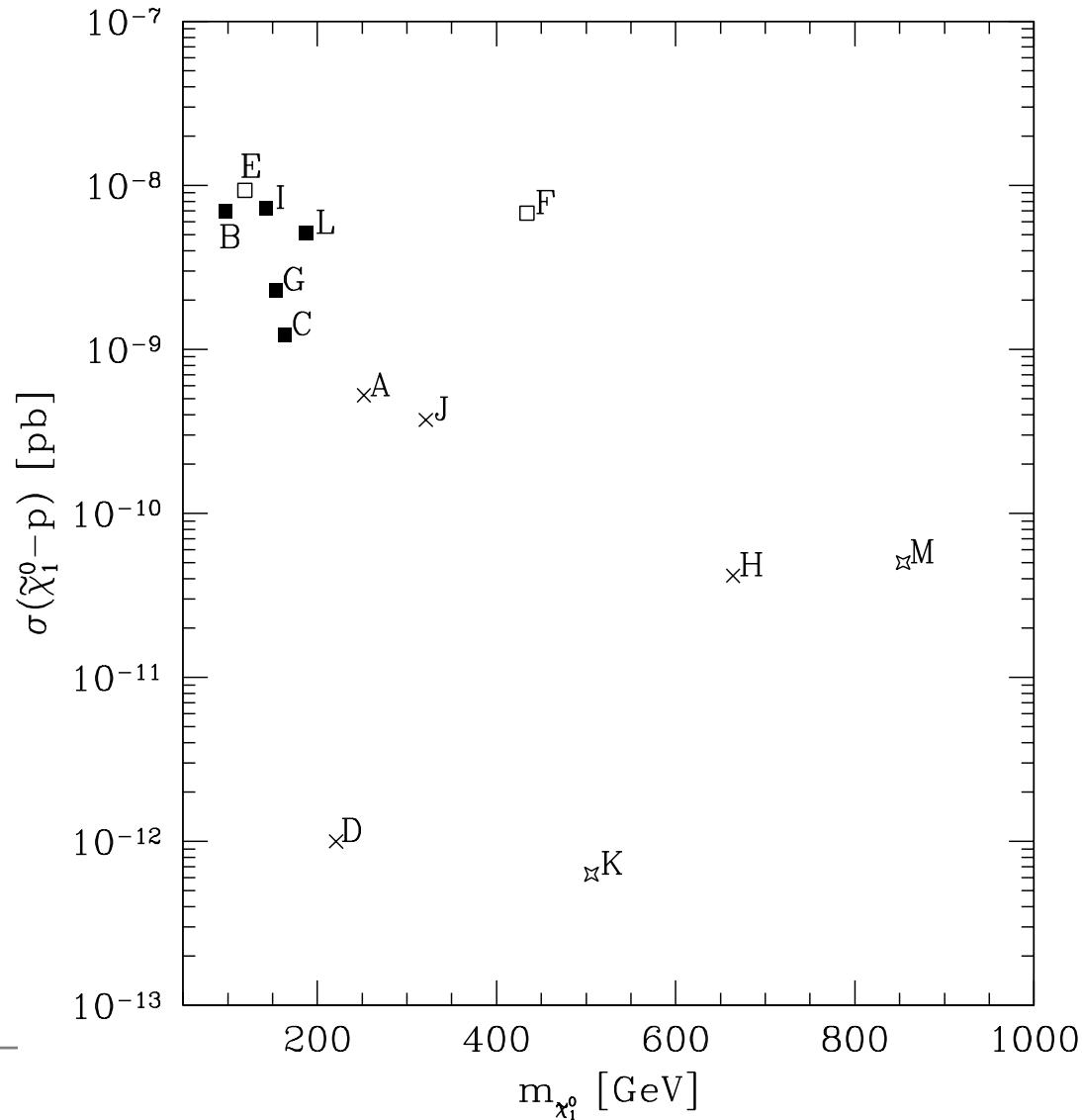
Points from Battaglia et al. (2003)



Solid squares: Bulk region  
Open squares: focus point region  
Crosses: Co-ann. region  
Stars: Higgs funnel

# Survey of Benchmark Points

Points from Battaglia et al. (2003)



Solid squares: Bulk region

Open squares: focus point region

Crosses: Co-ann. region

Stars: Higgs funnel

Consider A, E, G!

# Effect of Varying SUSY Parameter

Let's vary one (weak-scale) parameter by 20%, and compute the resulting change of  $\sigma_{\tilde{\chi}_1^0 p}$ !

# Effect of Varying SUSY Parameter

Let's vary one (weak-scale) parameter by 20%, and compute the resulting change of  $\sigma_{\tilde{\chi}_1^0 p}$ !

Point	$\sigma_{\tilde{\chi}_1^0}$ [pb]	$\delta\sigma(m_{\tilde{q}})$	$\delta\sigma(\mu)$	$\delta\sigma(\tan\beta)$	$\delta\sigma(m_A)$
A	$0.49 \times 10^{-9}$	-1.7%	-45.3%	-15.8%	-4.7%
E	$18.6 \times 10^{-9}$	-6.3%	-60.3%	-8.5%	-2.9%
G	$2.54 \times 10^{-9}$	-4.7%	-44.5%	+18%	-28%

# Lessons

- Relatively easily measurable squark mass has little influence: Higgs exchange dominates!



# Lessons

- Relatively easily measurable squark mass has little influence: Higgs exchange dominates!
- $\sigma_{\tilde{\chi}_1^0 p}$  is most sensitive to  $\mu$ , which determines gaugino–higgsino mixing: difficult to measure, except maybe in focus point region

# Lessons

- Relatively easily measurable squark mass has little influence: Higgs exchange dominates!
- $\sigma_{\tilde{\chi}_1^0 p}$  is most sensitive to  $\mu$ , which determines gaugino–higgsino mixing: difficult to measure, except maybe in focus point region
- If  $\tan \beta \gg 1$  (point G):  $\sigma_{\tilde{\chi}_1^0 p} \propto \tan^2 \beta / m_H^4$ : need parameters of Higgs sector!

# Summary

- WIMPs are still great CDM candidates!

# Summary

- WIMPs are still great CDM candidates!
- Learning from direct detection:

# Summary

- WIMPs are still great CDM candidates!
- Learning from direct detection:
  - Direct reconstruction of  $f_1(v)$  needs several hundred events

# Summary

- WIMPs are still great CDM candidates!
- Learning from direct detection:
  - Direct reconstruction of  $f_1(v)$  needs several hundred events
  - Non-trivial statements about moments of  $f_1$  possible with few dozen events

# Summary

- WIMPs are still great CDM candidates!
- Learning from direct detection:
  - Direct reconstruction of  $f_1(v)$  needs several hundred events
  - Non-trivial statements about moments of  $f_1$  possible with few dozen events
  - With  $\geq 2$  experiments: can get  $m_\chi$ !

# Summary

- WIMPs are still great CDM candidates!
- Learning from direct detection:
  - Direct reconstruction of  $f_1(v)$  needs several hundred events
  - Non-trivial statements about moments of  $f_1$  possible with few dozen events
  - With  $\geq 2$  experiments: can get  $m_\chi$ !
- Learning about direct detection ( $\sigma_{\tilde{\chi}_1^0 p}$ ):



# Summary

- WIMPs are still great CDM candidates!
- Learning from direct detection:
  - Direct reconstruction of  $f_1(v)$  needs several hundred events
  - Non-trivial statements about moments of  $f_1$  possible with few dozen events
  - With  $\geq 2$  experiments: can get  $m_\chi$ !
- Learning about direct detection ( $\sigma_{\tilde{\chi}_1^0 p}$ ):
  - Large hadronic uncertainty, especially from  $m_s \langle p | \bar{s}s | p \rangle$  (almost equivalently,  $\Sigma_{\pi N}$ )

# Summary

- WIMPs are still great CDM candidates!
- Learning from direct detection:
  - Direct reconstruction of  $f_1(v)$  needs several hundred events
  - Non-trivial statements about moments of  $f_1$  possible with few dozen events
  - With  $\geq 2$  experiments: can get  $m_\chi$ !
- Learning about direct detection ( $\sigma_{\tilde{\chi}_1^0 p}$ ):
  - Large hadronic uncertainty, especially from  $m_s \langle p | \bar{s}s | p \rangle$  (almost equivalently,  $\Sigma_{\pi N}$ )
  - In SUSY: Have to measure parameters of Higgs(ino) sector! Probably difficult to do at LHC.

# Summary

- WIMPs are still great CDM candidates!
- Learning from direct detection:
  - Direct reconstruction of  $f_1(v)$  needs several hundred events
  - Non-trivial statements about moments of  $f_1$  possible with few dozen events
  - With  $\geq 2$  experiments: can get  $m_\chi$ !
- Learning about direct detection ( $\sigma_{\tilde{\chi}_1^0 p}$ ):
  - Large hadronic uncertainty, especially from  $m_s \langle p | \bar{s}s | p \rangle$  (almost equivalently,  $\Sigma_{\pi N}$ )
  - In SUSY: Have to measure parameters of Higgs(ino) sector! Probably difficult to do at LHC.
- Both  $f_1(v)$  and  $\sigma_{\chi p}$  are needed to determine  $\rho_\chi$ : required input for learning about early Universe!