

Fixing D7 Brane Positions by F-Theory Fluxes

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A. Braun, A. Hebecker, CL, R. Valandro,

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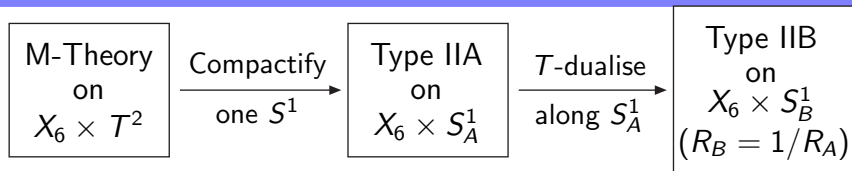
Motivation

- F-Theory: Nonperturbative version of type IIB string theory
[Vafa;Sen]
- Add two auxiliary dimensions, singularities of compactification manifold encode brane positions
- Recently, lots of interest in F-theory for (GUT) model building
[Beasley,Heckman,Vafa; Donagi,Wijnholt; Marsano,Saulina,Schäfer-Nameki; Bourjaily; Tatar,Watari,Hayashi,Kawanao,Tsuchiya; ...]
- Local models do not address global constraints like tadpole cancellation
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- Four-form flux can stabilise moduli, including brane positions
- Simple example: F-Theory on $K3 \times \widetilde{K3}$, where $\widetilde{K3}$ is an elliptic fibration over \mathbb{P}^1 [Görllich et al.; Lust et al.; Aspinwall,Kalosh;Dasgupta et al.]
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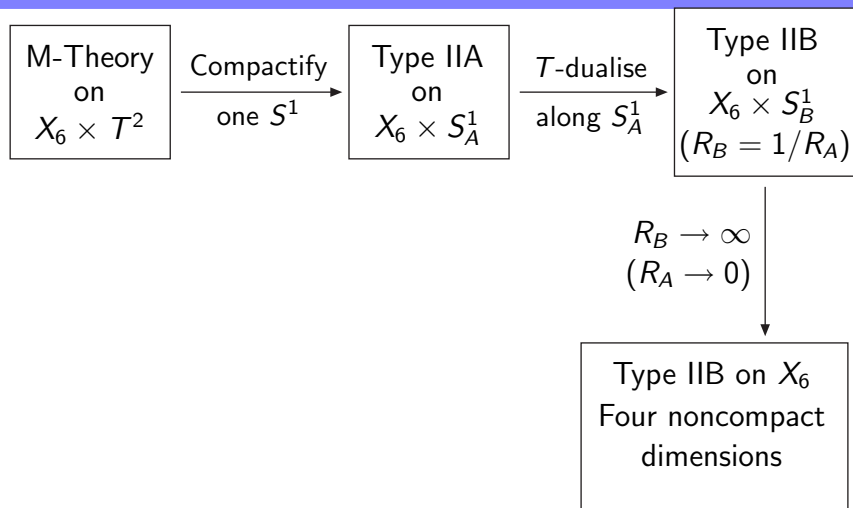
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F-Theory/M-Theory Duality



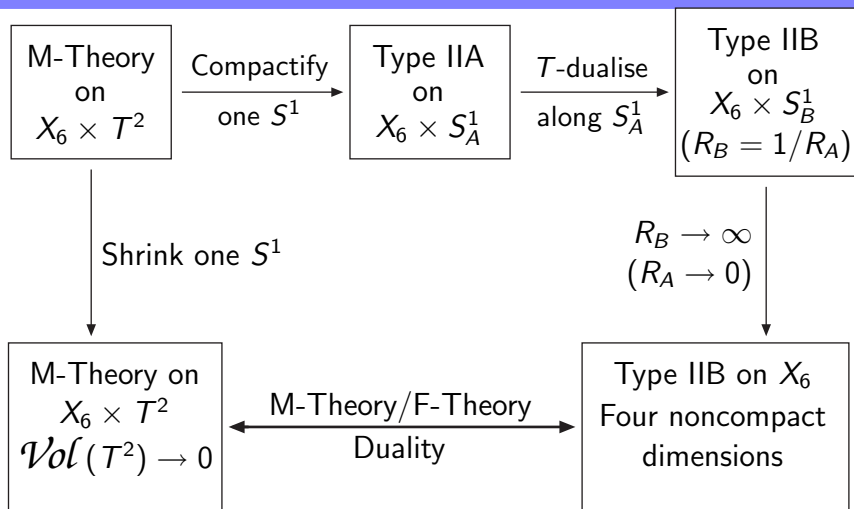
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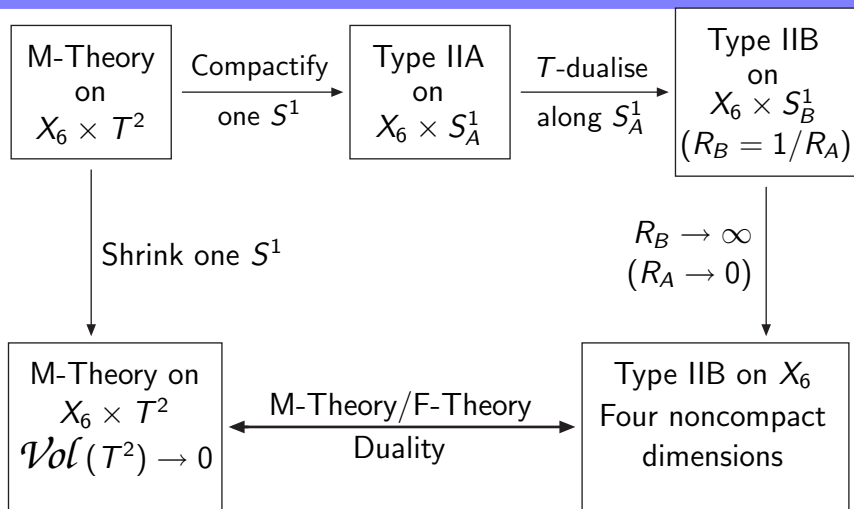
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K3: Calabi–Yau Two-Fold

- $H^2(K3, \mathbb{R})$ has signature $(3, 19)$
- Holomorphic two-form and Kähler form spanned by three real forms ω_i with $\omega_i \cdot \omega_j = \delta_{ij}$ and overall volume ν :

$$\omega = \omega_1 + i\omega_2 \quad j = \sqrt{2\nu} \omega_3$$

- K3 is **hyperkähler**: $SO(3)$ symmetry rotating the ω_i
 \rightsquigarrow geometry fixed by positive-norm **three-plane** $\Sigma \subset H^2(K3, \mathbb{R})$
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- Moduli space: Rotation of the ω_i orthogonal to Σ and volume
 \rightsquigarrow dimension = $3 \times 19 + 1 = 58$

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Cycles and Intersection Form

- There is an **integral basis** for $H^2(K3)$ such that intersection matrix is

$$U \oplus U \oplus U \oplus (-E_8) \oplus (-E_8) ,$$

where $U = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ and E_8 is Cartan matrix of E_8

- The “ U blocks” contain the positive norm directions
- ⇒ The ω_i must have components along the **U blocks**
Components along “ **E_8 directions**” determine gauge group

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K3: Elliptic Fibration and F-Theory Limit

- For an elliptically fibred $K3$, require integral cycles B and F (base and fibre) with

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- $B \cdot \omega = F \cdot \omega = 0$

$\Rightarrow (B + F, F)$ spans a U block, and we can parametrise the Kähler form as

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- F-theory limit: Fibre volume shrinks to zero $\Rightarrow b \rightarrow 0$. $K3$ volume is $\nu \sim bf - c^a c^a$, so we have to take $c^a \rightarrow 0$ as fast as \sqrt{b} (as intuitively expected)
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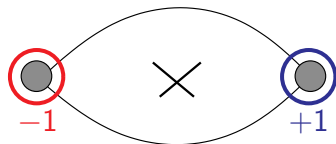
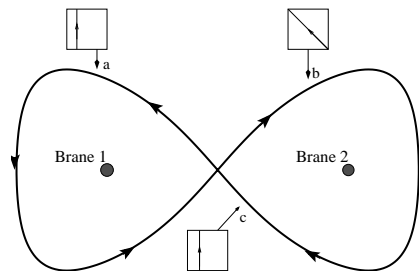
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Cycles Between Branes

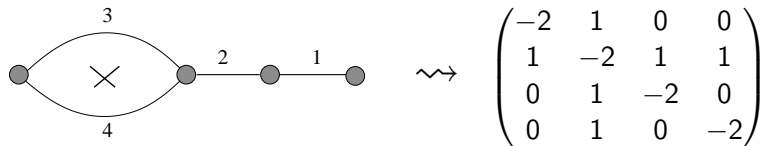
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- One leg in the base, one in the fibre torus
- **Shrink to zero** when the branes are moved on top of each other.
- They are topologically a **sphere** \leftrightarrow self-intersection -2 .
- Cycles meeting at a brane intersect once, cycles encircling 0 planes (\times) do not intersect

Shrinking Cycles and Gauge Enhancement

Intersection matrix of shrinking cycles determines gauge group:
Consider e.g. T^2/\mathbb{Z}_2 orientifold: One O7, four D7s $\rightsquigarrow SO(8)$



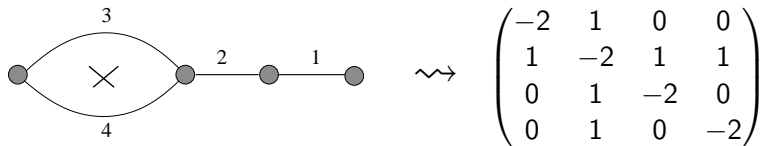
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$$\omega = \frac{\alpha}{2} + u e_2 + s \frac{\beta}{2} - \left(u s - \frac{z^2}{2} \right) e_1 + z_I \hat{E}_I$$

Explicit mapping between complex structure and brane positions!

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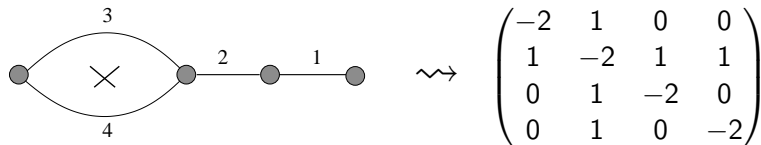
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base complex structure

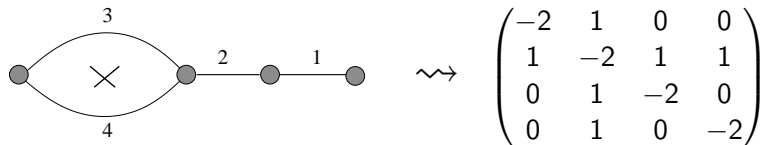
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- Type IIB: **Three-form flux G_3** on the bulk, **two-form gauge flux F_2** on the branes can stabilise geometric and brane moduli
- In M-theory, these are combined into **four-form flux G_4** (brane moduli become four-fold geometric moduli)
- Consistency conditions:
 - Flux quantisation: flux needs to be integral
 - Tadpole cancellation (without spacetime-filling M2 branes)

$$\frac{1}{2} \int_{K3 \times \widetilde{K3}} G_4 \wedge G_4 = \frac{\chi}{24} = 24$$

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$$V = \frac{1}{4\mathcal{V}^3} \left(\int_{K3 \times \widetilde{K3}} G_4 \wedge *G_4 - \frac{\chi}{12} \right)$$

- $K3 \times \widetilde{K3}$ is not a proper CY_4 : Holonomy is $SU(2) \times SU(2)$
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$K3 \times \widetilde{K3}$ Flux Potential

$$V = -\frac{1}{2(\nu \cdot \widetilde{\nu})^3} \left(\sum_j \|G \widetilde{\omega}_j\|_{\perp}^2 + \sum_i \|G^a \omega_i\|_{\perp}^2 \right)$$

$\|\cdot\|_{\perp}^2$ is the norm orthogonal to the three-plane

- Positive definite potential
- Manifestly symmetric under $SO(3)$
- Minima at $V = 0$:

$$G \widetilde{\omega}_j \in \langle \omega_1, \omega_2, \omega_3 \rangle \quad G^a \omega_i \in \langle \widetilde{\omega}_1, \widetilde{\omega}_2, \widetilde{\omega}_3 \rangle$$

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Minima: Existence, Flat Directions

- Minkowski minima do **not necessarily** exist: $G^a G$ must be diagonalisable and positive semi-definite (not guaranteed although $G^a G$ is self-adjoint, since metric is indefinite!)
- Flat directions **generally exist** and are **desired**: M-theory moduli become part of 4D **vector fields** in F-theory limit \rightsquigarrow fixing these moduli breaks the gauge group (rank-reducing)
- Flux also induces explicit mass term for three-dimensional vectors
- Vacua can preserve $\mathcal{N} = 4$, $\mathcal{N} = 2$ or $\mathcal{N} = 0$ supersymmetry in four dimensions, depending on the action of G on the three-plane

Stabilisation Strategy

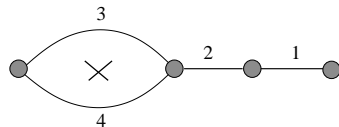
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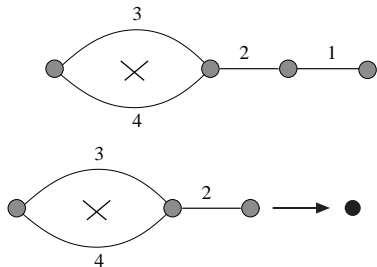
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- The T^2/\mathbb{Z}_2 orientifold with $SO(8)^4$: Four stacks of four D7 branes and one O7 plane each
- Moving one brane off a stack.
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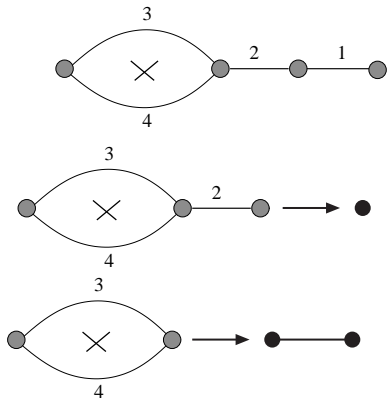
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